

# SUPERSYMMETRY BREAKING IN THE NO-SCALE SECTOR OF CALABI-YAU STRING MODELS

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- Soft scalar masses in supergravity
- Sequestering and global symmetries
- Generic Calabi-Yau string models
- Kähler potential and contact terms
- Structure of scalar masses

# SOFT SCALAR MASSES IN SUPERGRAVITY

## General structure of soft scalar masses

Kaplunovsky, Louis, 1993  
Brignole, Ibanez, Munoz 1993

In a supergravity theory with Kähler potential  $K$  and superpotential  $W$ , the soft scalar masses induced for the visible superfields  $Q^\alpha$  when the hidden superfields  $X^i$  break supersymmetry are given by:

$$m_{\alpha\bar{\beta}}^2 = - \left( R_{\alpha\bar{\beta}i\bar{j}} - \frac{1}{3} g_{\alpha\bar{\beta}} g_{i\bar{j}} \right) F^i \bar{F}^{\bar{j}}$$

This can be written in terms of the Kähler function  $\Omega = -3 e^{-K/3}$  as:

$$\begin{aligned} m_{\alpha\bar{\beta}}^2 &= 3 \Omega^{-1} \left( \Omega_{\alpha\bar{\beta}i\bar{j}} - \Omega^{-1} \bar{\delta}^{\gamma} \Omega_{\bar{\delta}\alpha i} \Omega_{\gamma\bar{\beta}\bar{j}} \right) F^i \bar{F}^{\bar{j}} \\ &= 3 \Omega^{-1} \left( \Omega_{\alpha\bar{\beta}} \Big|_D - \Omega^{-1} \bar{\delta}^{\gamma} \Big| \Omega_{\bar{\delta}\alpha} \Big|_F \Omega_{\gamma\bar{\beta}} \Big|_{\bar{F}} \right) \end{aligned}$$

The crucial ingredients are thus the operators in  $\Omega$  that mix  $Q^\alpha$  and  $X^i$ , and the orientation of the Goldstino direction.

# THE SUPERSYMMETRIC FLAVOR PROBLEM

## Flavor structure of soft scalar masses

The **flavor structure** of the soft scalar mass matrix  $m_{\alpha\bar{\beta}}^2$  is a priori **generic**, because this is generated at the fundamental scale of the theory where the **flavor structure** of the ordinary fermion masses must also emerge.

This causes severe problems at the phenomenological level. One should then find some mechanism that naturally forces  $m_{\alpha\bar{\beta}}^2$  to be approximately **flavor-universal**.

## Possible explanations for universality

The two most interesting ideas to explain **flavor-universality** of soft scalar masses are **mode sequestering** along extra dimensions and **selection rules** from global symmetries.

# SEQUESTERING ALONG EXTRA DIMENSIONS

## Localization along an extra dimension

Randall, Sundrum 1999

Anisimov, Dine, Graesser, Thomas 2002

If the **visible** and **hidden** sectors are separated along an extra dimension, as it happens quite naturally in string models, the classical contribution to  $m_{\alpha\bar{\beta}}^2$  is very restricted by locality.

In the **minimal** case with only the **gravity** multiplet in the bulk  $m_{\alpha\bar{\beta}}^2 = 0$ . In the **general** case involving also **vector** multiplets in the bulk  $m_{\alpha\bar{\beta}}^2 \neq 0$ , but its form is very special and one may force it to vanish by other means.

When  $m_{\alpha\bar{\beta}}^2 = 0$  at the classical level, one may get a viable  $m_{\alpha\bar{\beta}}^2 \neq 0$  from quantum corrections, which are approximately flavor-universal and insensitive to high-energy physics.

It would be interesting to understand whether and how this situation can be implemented in generic string models.

# HIDDEN GLOBAL SYMMETRIES

## Constraints from global symmetries

Schmaltz, Sundrum 2006

Andrey, Scrucra 2010

If the hidden sector possesses global symmetries generated by some Killing vectors  $k_a^i$  with potentials  $D_a$ , the form of  $m_{\alpha\bar{\beta}}^2$  gets constrained. Most importantly, the Goldstino direction is restricted to satisfy:

$$\bar{k}_{a i} F^i = -i D_a m_{3/2} \quad \nabla_i k_{a \bar{j}} F^i \bar{F}^{\bar{j}} = -2i D_a m_{3/2}^2$$

In situations where  $D_a \simeq 0$  on the vacuum, the gravity effects on the right-hand sides trivialize and one then finds:

$$\bar{k}_{a i} F^i \simeq 0 \quad \nabla_i k_{a \bar{j}} F^i \bar{F}^{\bar{j}} \simeq 0$$

In rigid superspace, these follow from the conservation law  $\mathcal{D}^2 J_a \simeq 0$  for the Nöther current  $J_a \simeq \text{Im}(K_i k_a^i)$ , and read  $J_a|_F \simeq 0$  and  $J_a|_D \simeq 0$ . It follows that contact terms in  $\Omega$  that involve  $J_a$  do not contribute to  $m_{\alpha\bar{\beta}}^2$ .

# MILD SEQUESTERING BY GLOBAL SYMMETRIES

## Contact terms and global symmetries

Kachru, McAllister, Sundrum 2007

For generic sequestered models, we expect  $\Omega$  to have the form

$$\Omega \simeq Q^\alpha \bar{Q}^{\bar{\alpha}} + X^i \bar{X}^{\bar{i}} + \frac{1}{2} M^{-2} \sum_a \left( c_{\alpha\bar{\beta}}^a Q^\alpha \bar{Q}^{\bar{\beta}} + c_{i\bar{j}}^a X^i \bar{X}^{\bar{j}} \right)^2$$

The soft scalar masses induced by a generic  $W$  are then given by

$$m_{\alpha\bar{\beta}}^2 \simeq M^{-2} \left[ - c_{\alpha\bar{\beta}}^a c_{i\bar{j}}^a F^i \bar{F}^{\bar{j}} + (c^a c^b)_{\alpha\bar{\beta}} c_{i\bar{j}}^a F^i X^{\bar{j}} c_{p\bar{q}}^b X^p \bar{F}^{\bar{q}} \right]$$

In presence of global symmetries with Nöther currents  $J_X^a \simeq c_{i\bar{j}}^a X^i \bar{X}^{\bar{j}}$ , the Ward identities  $J_X^a|_F \simeq 0$  and  $J_X^a|_D \simeq 0$  imply:

$$c_{i\bar{j}}^a F^i \bar{X}^{\bar{j}} \simeq 0 \quad c_{i\bar{j}}^a F^i \bar{F}^{\bar{j}} \simeq 0$$

In such a situation one then finds:

$$m_{\alpha\bar{\beta}}^2 \simeq 0$$

# HETEROTIC STRING MODELS

## Heterotic M-theory on a Calabi-Yau

Horava, Witten 1996

Lukas, Ovrut, Stelle, Waldram 1999

Let us consider a generic heterotic string model based on a Calabi-Yau manifold  $M$  and a stable holomorphic vector bundle  $E_v \times E_h$  over it. This also arises from M-theory on  $M \times S^1/Z_2$  with two sequestered branes, in the weakly coupled limit where the size of  $S^1/Z_2$  is small.

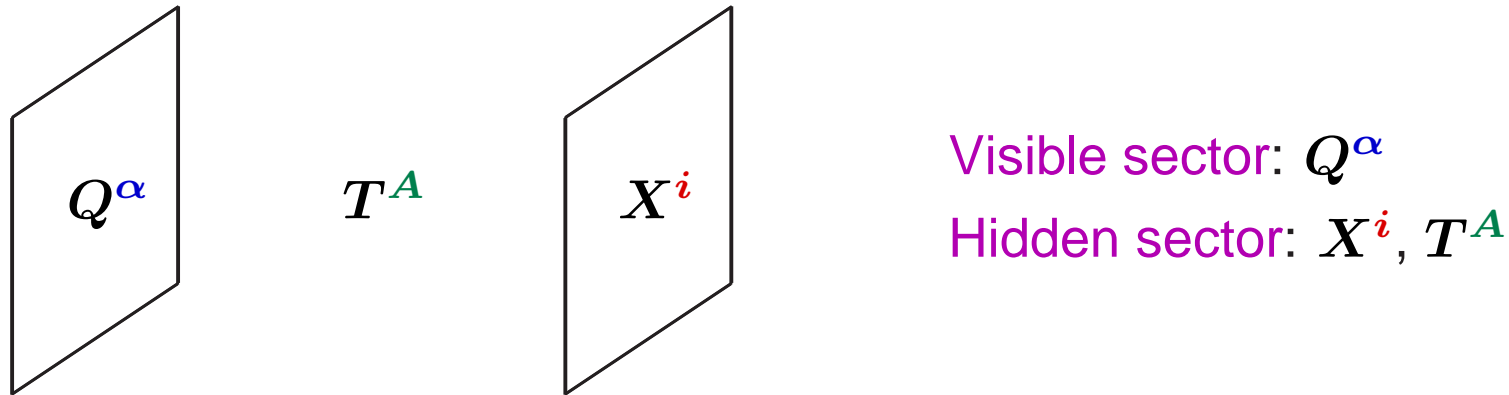
The 4D effective theory can be lifted to a 5D theory with localized brane sectors containing visible and hidden matter superfields  $Q^\alpha$  and  $X^i$ , and a bulk sector containing in particular Kähler moduli superfields  $T^A$ .

The non-minimal Kähler moduli  $T^a$  come along with heavy vectors  $V^a$ , which couple non-trivially to  $Q^\alpha$ ,  $X^i$  and  $T^A$ . When integrated out, these induce contact terms in the effective Kähler function  $\Omega$ .

# BRANE WORLD INTERPRETATION

## Geometric picture

The **brane** sectors contain only the light modes  $Q^\alpha$  and  $X^i$ , while the **bulk** sector contains the light modes  $T^A$  plus heavy Kaluza-Klein modes.



The masses  $m_{\alpha\bar{\beta}}^2$  get two kinds of contributions: one from operators in  $\Omega$  mixing  $Q^\alpha$  and  $X^i$ , and one from operators in  $\Omega$  mixing  $Q^\alpha$  and  $T^A$ . We thus need to know the full dependence of  $K$  on  $Q^\alpha$ ,  $X^i$  and  $T^A$ .



# DERIVATION OF THE EFFECTIVE THEORY

## Reduction of the standard heterotic string

Witten 1985

The light 4D fields arise from the possible zero-modes of the 10D fields. The  $Q^\alpha$ ,  $X^i$  come from harmonic 1-forms in  $H^1(M, E_v)$ ,  $H^1(M, E_h)$ , while the  $T^A$  come from harmonic (1, 1)-forms in  $H^{1,1}(M)$ :

$$Q^\alpha \Leftrightarrow u_\alpha \quad X^i \Leftrightarrow u_i \quad T^A \Leftrightarrow \omega_A$$

The effective  $K$  for the light fields may be derived by working out their kinetic terms by reduction on  $M$  and comparing with the general structure of supergravity theories. The effective  $W$  can be imagined to be arbitrary.

Discarding rather than integrating out heavy non-zero modes associated to non-harmonic forms is justified only whenever:

$$\text{tr}(u_\alpha \wedge \bar{u}_{\bar{\beta}}) \text{ and } \text{tr}(u_i \wedge \bar{u}_{\bar{j}}) \text{ harmonic} \Leftrightarrow \omega_A$$

# EFFECTIVE KÄHLER POTENTIAL

General result for matter fields and Kähler moduli

Candelas, de la Ossa 1990  
Paccetti Correia, Schmidt 2008  
Andrey, Scrucca 2011

The effective Kähler potential is found to be

$$K = -\log \left( d_{ABC} J^A J^B J^C \right)$$

where

$$J^A = T^A + \bar{T}^A - c_{\alpha\bar{\beta}}^A Q^\alpha \bar{Q}^{\bar{\beta}} - c_{i\bar{j}}^A X^i \bar{X}^{\bar{j}}$$

The numerical quantities defining this result are:

$$d_{ABC} = \int \omega_A \wedge \omega_B \wedge \omega_C$$
$$c_{\alpha\bar{\beta}}^A = \int \omega^A \wedge \text{tr}(u_\alpha \wedge \bar{u}_{\bar{\beta}}) \quad c_{i\bar{j}}^A = \int \omega^A \wedge \text{tr}(u_i \wedge \bar{u}_{\bar{j}})$$

This extends the results for the special case of the untwisted sector of orbifolds, where harmonic forms are covariantly constant, to a larger class of cases, where harmonic forms close under multiplication.

# CANONICAL PARAMETRIZATION

## Canonical basis

With a suitable basis for the harmonic forms and fields, which defines an overall modulus  $T$  and some relative moduli  $T^a$ , one may rewrite  $K$  as:

$$K = -\log \left( J^3 - \frac{1}{2} J J^a J^a + \frac{1}{6} d_{abc} J^a J^b J^c \right)$$

where

$$J = T + \bar{T} - \frac{1}{3} Q^\alpha \bar{Q}^{\bar{\alpha}} - \frac{1}{3} X^i \bar{X}^{\bar{i}}$$
$$J^a = T^a + \bar{T}^a - c_{\alpha\bar{\beta}}^a Q^\alpha \bar{Q}^{\bar{\beta}} - c_{i\bar{j}}^a X^i \bar{X}^{\bar{j}}$$

## Contact terms

The leading terms in the Kähler function for  $J^a \ll J$  are

$$\Omega \simeq -3J + \frac{1}{2} J^{-1} J^a J^a - \frac{1}{6} d_{abc} J^{-2} J^a J^b J^c$$

# STRUCTURE OF SOFT SCALAR MASSES

## Structure of soft terms

One may study the following reference point, around which the canonical parametrization is particularly convenient:

$$T \simeq \frac{1}{2} \quad T^a \simeq 0 \quad Q^\alpha \simeq 0 \quad X^i \simeq 0$$

Using the computed  $K$  and imagining a generic  $W$ , one obtains:

$$m_{\alpha\bar{\beta}}^2 \simeq -c_{\alpha\bar{\beta}}^a c_{i\bar{j}}^a F^i \bar{F}^{\bar{j}} - \left( \frac{1}{3} \delta_{\alpha\bar{\beta}} \delta_{ab} + (d_{abc} c^c - c^a c^b)_{\alpha\bar{\beta}} \right) F^a \bar{F}^{\bar{b}} \\ - c_{\alpha\bar{\beta}}^a F^a \bar{F} + \text{c.c.}$$

This vanishes identically if the Goldstino direction is suitably constrained:

$$m_{\alpha\bar{\beta}}^2 \simeq 0 \Leftrightarrow F^a \simeq 0 \quad \text{and} \quad c_{i\bar{j}}^a F^i \bar{F}^{\bar{j}} \simeq 0$$

# SEQUESTERING BY GLOBAL SYMMETRIES

## Approximate global symmetries

Andrey, Scrucca 2011

The Goldstino direction can be guaranteed to point in a direction for which  $m_{\alpha\bar{\beta}}^2 \simeq 0$  by postulating that the following transformations represent two approximate symmetries not only of  $K$  but also of  $W$ :

$$\begin{aligned}\delta_a^1 T^b &= i\delta_a^b \Leftrightarrow F^a \simeq 0 \\ \delta_a^2 X^i &= -ic_{\bar{j}i}^a X^{\bar{j}} \Leftrightarrow c_{i\bar{j}}^a F^i \bar{F}^{\bar{j}} \simeq 0\end{aligned}$$

Clearly  $\delta_a^1$  always form a group  $R^n$  and give exact symmetries of  $K$ . However  $\delta_a^2$  only form a group  $H$  if  $c_{i\bar{j}}^a$  generate a closed algebra and only extends to exact symmetries of  $K$  if  $d_{abc}$  is a symmetric invariant of this algebra. We conclude that:

Sequestering by symmetries possible for some Calabi-Yau models

# PARTICULAR CASE OF ORBIFOLDS

Ferrara, Kounnas, Porrati 1986

Li, Peschanski, Savoy 1987

Andrey, Scrucca 2010

## Symmetric structure in orbifolds

One special class of models where one is automatically in business is provided by orbifold constructions. In the untwisted sector, the formula for  $K$  that has been obtained applies, with:

$c_{\alpha\bar{\beta}}^a, c_{i\bar{j}}^a$  : generators of some  $H \subset SU(3)$

$d_{abc}$  : symmetric invariant of this  $H \subset SU(3)$

The scalar manifold is always a symmetric coset manifold, and  $H$  belongs to the stability group. As a result,  $R^n \times H$  is an exact symmetry of  $K$ , and imposing it also to  $W$  leads to vanishing masses. We thus conclude that:

Sequestering by symmetries possible for all orbifold models

# ORIENTIFOLD STRING MODELS

## Low-energy effective theory

Grimm, Louis 2004

Grana, Grimm, Jockers, Louis 2004

Jockers, Louis 2004

In orientifold models, one can achieve a similar setup. The low-energy effective theory is technically more involved but qualitatively similar to that of heterotic models. It is again parametrized by the intersection numbers  $d_{abc}$  and some matrices  $c_{\alpha\bar{\beta}}^a$  and  $c_{i\bar{j}}^a$ .

## Soft scalar masses

Work in progress

The structure of soft scalar masses is similar to that of heterotic models and mild sequestering seems to be possible under similar circumstances. Again, it appears to work in all models based on orbifolds, but only in a subclass of those based on smooth Calabi-Yaus.

# CONCLUSIONS

- Under some assumptions, the Kähler potential of the no-scale sector of string models can be fully computed. The resulting soft scalar masses do vanish for suitably oriented Goldstino directions.
- The Goldstino direction can be forced to align along such special directions by relying on some global symmetries, but this appears to be possible only under some extra assumptions.
- A special class of models where this mechanism can always work is that of orbifold models. But it might be possible to put it at work also for other special classes of Calabi-Yau models.