

# ÉCOLE POLYTECHNIQUE FÉDÉRALE DE LAUSANNE

School of Computer and Communication Sciences

**Handout 4**  
Homework 4

Statistical Physics for Communication and Computer Science  
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In homework 2 you proved that the Ising model in one dimension ( $d = 1$ ) does not have a phase transition for any  $T > 0$ . On the grid  $\mathbb{Z}^d$  there is a non trivial phase diagram with first and second order phase transitions for any  $d \geq 2$ . This is also the case on the complete graph (as shown in the lectures) which morally corresponds to  $d = +\infty$ . Another graph that in a sense, corresponds to  $d = +\infty$ , is the  $q$ -ary tree for  $q \geq 3$ . Indeed on  $\mathbb{Z}^d$  the number of lattice sites at distance less than  $n$  from the origin scales as  $n^d$ . On the  $q$ -ary tree it scales as  $(q - 1)^n$  which grows faster than  $n^d$  for any finite  $d$  (for  $q \geq 3$ ). Of course  $q = 2$  corresponds to  $\mathbb{Z}_+$ .

The goal of the two exercises below is to solve for the Ising model on a  $q$ -ary tree and show that it displays first and second order phase transitions (with similar qualitative properties than on a complete graph).

Consider a finite rooted tree and call the root vertex  $o$ . All vertices have degree  $q$ , except for the leaf nodes that have degree 1. We suppose that the tree has  $n$  levels (the root being “level 0”). The thermodynamic limit corresponds to  $n \rightarrow +\infty$ . The Hamiltonian (multiplied by  $\beta$ ) is

$$\beta\mathcal{H}_n = -K \sum_{(i,j) \in E_n} s_i s_j - h \sum_{i \in V_n} s_i \quad (1)$$

where  $K > 0$ ,  $h \in \mathbb{R}$ ,  $V_n$  is the set of vertices and  $E_n$  the set of edges. We are interested in the magnetization of the root node in the thermodynamic limit:

$$m(K, h) = \lim_{n \rightarrow +\infty} \langle s_o \rangle_n = \frac{\sum_{\{s_k, k \in V_n\}} s_o e^{-\beta\mathcal{H}_n}}{Z_n} \quad (2)$$

The formula  $\tanh^{-1} y = \frac{1}{2} \ln \frac{1+y}{1-y}$  might be useful.

**Problem 1** (Recursive equations). Perform the sums over the spins attached at the leaf nodes and show that

$$\langle s_o \rangle_n = \frac{\sum_{\{s_k, k \in V_{n-1}\}} s_o e^{-\beta\mathcal{H}'_{n-1}}}{Z'_{n-1}} \quad (3)$$

where  $E_{n-1}$  and  $V_{n-1}$  are the edge and vertex sets of a tree with  $n - 1$  levels and the new Hamiltonian is

$$\beta\mathcal{H}'_n = -K \sum_{(i,j) \in E_{n-1}} s_i s_j - h \sum_{i \in V_{n-1}} s_i - (q - 1) \tanh^{-1}(\tanh K \tanh h) \sum_{i \in \text{level } n-1} s_i \quad (4)$$

Iterate this calculation and deduce

$$\langle s_o \rangle_n = \tanh(h + q \tanh^{-1}(\tanh K \tanh u_n)) \quad (5)$$

where

$$u_{k+1} = h + (q - 1) \tanh^{-1}(\tanh K \tanh u_k), \quad u_1 = h \quad (6)$$

Check that for  $q = 2$  you get back the recursion of homework 2.

**Problem 2** (Analysis of the recursion). We want to analyze the fixed point equation for  $q \geq 3$ ,

$$u = h + (q - 1) \tanh^{-1}(\tanh K \tanh u) \quad (7)$$

Plot the curves  $u \rightarrow u - h$  and  $u \rightarrow (q - 1) \tanh^{-1}(\tanh K \tanh u)$  and show that:

- for  $K \leq K_c \equiv \frac{1}{2} \ln \frac{q}{q-2} = \tanh^{-1}(q - 1)^{-1}$ , (7) has a unique solution, and that the iterations (6) converge to this unique solution.
- for  $K > K_c$ :
  - for  $|h| \geq h_s$ , (7) has a unique solution (you do not needw3 to compute  $h_s$  explicitly although it is possible to find its analytical expression) and that the iterations (6) converge to this unique solution.
  - for  $|h| < h_s$ , (7) has three solutions  $u_-(h) < u_0(h) < u_+(h)$ . Check graphically that for  $h > 0$  the iterations (6) with initial condition  $u_1 = h$  converge to  $u_+(h)$ . Similarly for  $h < 0$  they converge to  $u_-(h)$ . Check also graphically that the fixed point  $u_0(h)$  is unstable whereas  $u_{\pm}(h)$  are stable.

**Problem 3** (Phase transitions). Now we want to discuss the consequences of the results in problem 2 for the phase diagram. On a tree the magnetization is defined as the average spin of the root

$$m(K, h) = \lim_{n \rightarrow +\infty} \langle s_o \rangle_n, \quad (8)$$

and we define the "spontaneous magnetization" as  $m_{\pm}(K) = \lim_{h \rightarrow 0_{\pm}} m(K, h)$ .

You will show that in the  $(K^{-1}, h)$  plane there is a first order phase transition line  $(K^{-1} \in [0, K_c^{-1}[, h = 0)$  terminated by a critical point  $K_c$ . Outside of this line  $m(K, h)$  is an analytic function of each variable.

- Deduce from the analysis in problem 2 that for  $K \leq K_c$ ,  $m_+(K) = m_-(K) = 0$ .
- Deduce that for  $K > K_c$ ,  $m_+(K) \neq m_-(K)$  (jump discontinuity or first order phase transition) and that for  $K \rightarrow +\infty$   $m_{\pm} \rightarrow \pm 1$ .
- Show that for  $K \rightarrow K_c$  from above,  $m_{\pm}(K) \sim (K - K_c)^{1/2}$ . So on the line  $h = 0$ , as a function of  $K$ , the spontaneous magnetization is continuous but not differentiable at  $K_c$  (second order phase transition).
- Now fix  $K = K_c$  and show that  $m(K_c, h) \sim |h|^{1/3}$ . As a function of  $h$  the spontaneous magnetization is continuous but not differentiable at  $K_c$  (second order phase transition).

**Hint:** for the last two questions you can expand the fixed point equation to order  $u^3$ .

**Remark 1:** Note that the exponents  $1/2$  and  $1/3$  are the same than for the model on a complete graph. This is also the case for all  $d \geq 4$  and is not the case for  $d = 2, 3$ .

**Remark 2:** On a tree the definition of the magnetization above is *not equivalent* to minus the derivative of the free energy with respect to  $h$ . In fact there is a fine point:  $-\frac{1}{n} \ln Z_n$  is dominated by the contributions of leaf nodes and is not the "physically meaningful" definition of free energy. Rather the "physically meaningful" definition is given by an integral with respect to  $h$  of the magnetization at the root