

3) End of lecture 11: Circuit and complexity of the QFT.

The QFT is defined by its action on basis vectors $\{|0\rangle, |1\rangle, \dots, |N-1\rangle\}$ of an N -dimensional Hilbert space $\mathcal{H} = \text{span}\{|0\rangle, \dots, |N-1\rangle\}$:

$$\text{QFT}|x\rangle = \frac{1}{\sqrt{N}} \sum_{y=0}^{N-1} \exp\left(2\pi i \frac{xy}{N}\right) |y\rangle.$$

a) Note that in the special case $N=2$ the QFT becomes the usual Hadamard gate:

$$\begin{aligned} (\text{QFT})_{N=2}|x\rangle &= \frac{1}{\sqrt{2}} \sum_{y=0}^1 \exp(\pi i xy) |y\rangle = \frac{1}{\sqrt{2}} \sum_{y=0}^1 (-1)^{xy} |y\rangle \\ &= \frac{1}{\sqrt{2}} (|0\rangle + (-1)^x |1\rangle) \end{aligned}$$

Of course the circuit in this case is simply



b) Let examine the case $N=4$ i.e. $\mathcal{H} = \text{span}\{|0\rangle, |1\rangle, |2\rangle, |3\rangle\}$.

$$\text{QFT}|x\rangle = \frac{1}{\sqrt{4}} \left\{ e^{2\pi i \frac{x \cdot 0}{4}} |0\rangle + e^{2\pi i \frac{x \cdot 1}{4}} |1\rangle + e^{2\pi i \frac{x \cdot 2}{4}} |2\rangle + e^{2\pi i \frac{x \cdot 3}{4}} |3\rangle \right\}.$$

We can represent the $|y\rangle$ in binary notation:

②

$$|0\rangle = |00\rangle ; |1\rangle = |01\rangle ; |2\rangle = |10\rangle ; |3\rangle = |11\rangle .$$

Then

$$\begin{aligned} \text{QFT}|x\rangle &= \frac{1}{\sqrt{4}} \left(|00\rangle + e^{i\frac{\pi}{2}x} |01\rangle + e^{i\pi x} |10\rangle + e^{3i\frac{\pi}{2}x} |11\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left(|0\rangle + e^{i\pi x} |2\rangle \right) \otimes \left(|0\rangle + e^{i\frac{\pi}{2}x} |1\rangle \right) . \end{aligned}$$

Now we can represent x also in binary notation:

$$x = \underbrace{2 \cdot x_1 + x_0}_{\substack{\{0; 1; 2; 3\} \\ m}} ; x_0, x_1 \in \{0, 1\} .$$

$$\Rightarrow e^{i\pi x} = e^{2\pi i x_1} e^{i\pi x_0} = (-1)^{x_0}$$

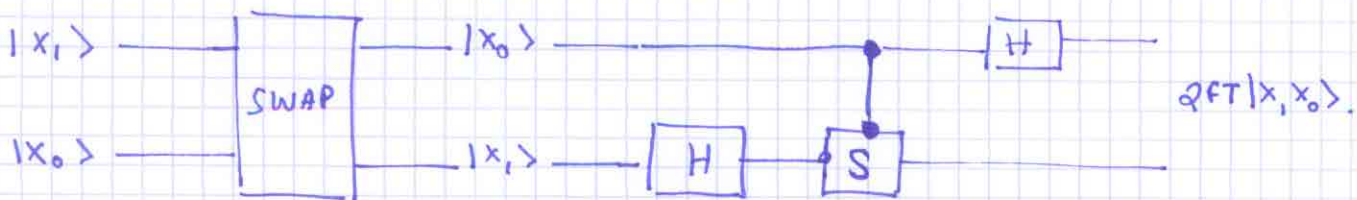
$$e^{i\frac{\pi}{2}x} = e^{i\pi x_1} e^{i\frac{\pi}{2}x_0} = (-1)^{x_1} e^{i\frac{\pi}{2}x_0} .$$

Thus

$$\text{QFT}|x\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle + (-1)^{x_0} |2\rangle \right) \otimes \left(|0\rangle + (-1)^{x_1} e^{i\frac{\pi}{2}x_0} |1\rangle \right)$$

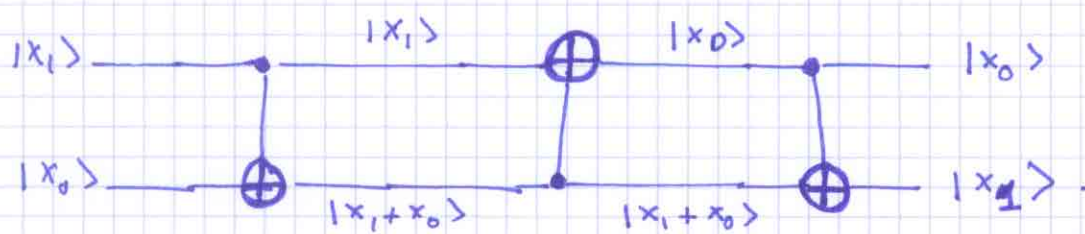
$$\text{QFT}|x_1, x_0\rangle .$$

A circuit realizing this operation is:



(3)

where the SWAP operation is realized as follows:



Once the SWAP operation is performed on $|x_1, x_0\rangle$ we act with H on the second qubit:

$$H \text{ SWAP } |x_1, x_0\rangle = H |x_0, x_1\rangle = |x_0\rangle \cdot \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} |1\rangle).$$

Then we act with a controlled $S = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/2} \end{pmatrix}$ gate:

$$\begin{aligned} CS H \text{ SWAP } |x_1, x_0\rangle &= CS |x_0\rangle \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} |1\rangle) \\ &= |x_0\rangle \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} e^{i\frac{\pi}{2} x_0} |1\rangle). \end{aligned}$$

Note that here the CS has two control bits: the first and the second. In fact the matrix for CS (a two bit gate) is

$$CS = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & e^{i\frac{\pi}{2}} \end{pmatrix}$$

so that $e^{i\pi/2}$ acts only on $|11\rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$. Another way to express CS is

$$CS = \underbrace{|0\rangle\langle 0|}_{\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}} \otimes \mathbb{1} + \underbrace{|1\rangle\langle 1|}_{\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}} \otimes S = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/2} \end{pmatrix}.$$

The last Hadamard gate acts on the first bit and yields:

$$\begin{aligned} H CS H \text{ SWAP } |x_1, x_0\rangle &= H |x_0\rangle \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} e^{i\frac{\pi}{2}x_0} |1\rangle) \\ &= \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_0} |1\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_1} e^{i\frac{\pi}{2}x_0} |1\rangle) \end{aligned}$$

So we have the decomposition (for $N=4$)

$$\text{qFT} = (H \otimes I) (CS) (I \otimes H) (\text{SWAP})$$

c) This decomposition and the corresponding circuit can easily be generalized to any $N = 2^m$.

Lemma: Let $x \in \{0, 1, \dots, N-1\}$ with $N = 2^m$.

$$\text{qFT} |x\rangle = \frac{1}{\sqrt{N}} \sum_{\ell=0}^{N-1} \left(|0\rangle + e^{i\frac{2\pi}{N}\ell x} |1\rangle \right)$$

Proof: Use the binary representation for $|y\rangle = |y_{m-1} \dots y_0\rangle$

where
$$y = 2^{m-1} y_{m-1} + 2^{m-2} y_{m-2} + \dots + 2^0 y_0$$

and the bits $y_i \in \{0, 1\}$. Take the definition of $\text{qFT} |x\rangle$

and split the sum over $y \in \{0, \dots, N-1\}$ into a sum over even

terms and a sum over odd terms;

(5)

$$\begin{aligned}
 \text{qFT } |x\rangle &= \frac{1}{2^{m/2}} \sum_{y \text{ even}} e^{2\pi i \frac{xy}{2^m}} |y\rangle + \sum_{y \text{ odd}} e^{2\pi i \frac{xy}{2^m}} |y\rangle \\
 &= \sum_{y'=0}^{2^{m-1}-1} e^{2\pi i \frac{x \cdot 2y'}{2^m}} |y_{m-1} \dots y_1 0\rangle + \sum_{y'=0}^{2^{m-1}-1} e^{2\pi i \frac{x \cdot (2y'+1)}{2^m}} |y_{m-1} \dots y_1 1\rangle
 \end{aligned}$$

where we used the facts that if $y = 2y'$ and

$$y = 2^{m-1} y_{m-1} + \dots + 2^1 y_1 + 2^0 y_0 \quad \text{then}$$

$$y' = 2^{m-2} y_{m-1} + \dots + 2^0 y_1 \quad \text{and } y_0 = 0.$$

if $y = 2y'+1$ and

$$y = 2^{m-1} y_{m-1} + \dots + 2^1 y_1 + 2^0 y_0 \quad \text{then}$$

$$y' = 2^{m-2} y_{m-1} + \dots + 2^0 y_1 \quad \text{and } y_0 = 1.$$

For this decomposition we conclude that

$$\text{qFT } |x\rangle = \left(\frac{1}{2^{(m-1)/2}} \sum_{y=0}^{2^{m-1}-1} e^{2\pi i \frac{xy}{2^{m-1}}} |y\rangle \right) \otimes \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i \frac{px}{2^{m-1}}} |1\rangle)$$

By repeating the same decomposition again and again on the first parenthesis we obtain the result of the Lemma. ■

The l -th term in the product (Lemma) is :

$$\left(|0\rangle + e^{i \frac{\pi}{2^{l-1}} x} |1\rangle \right).$$

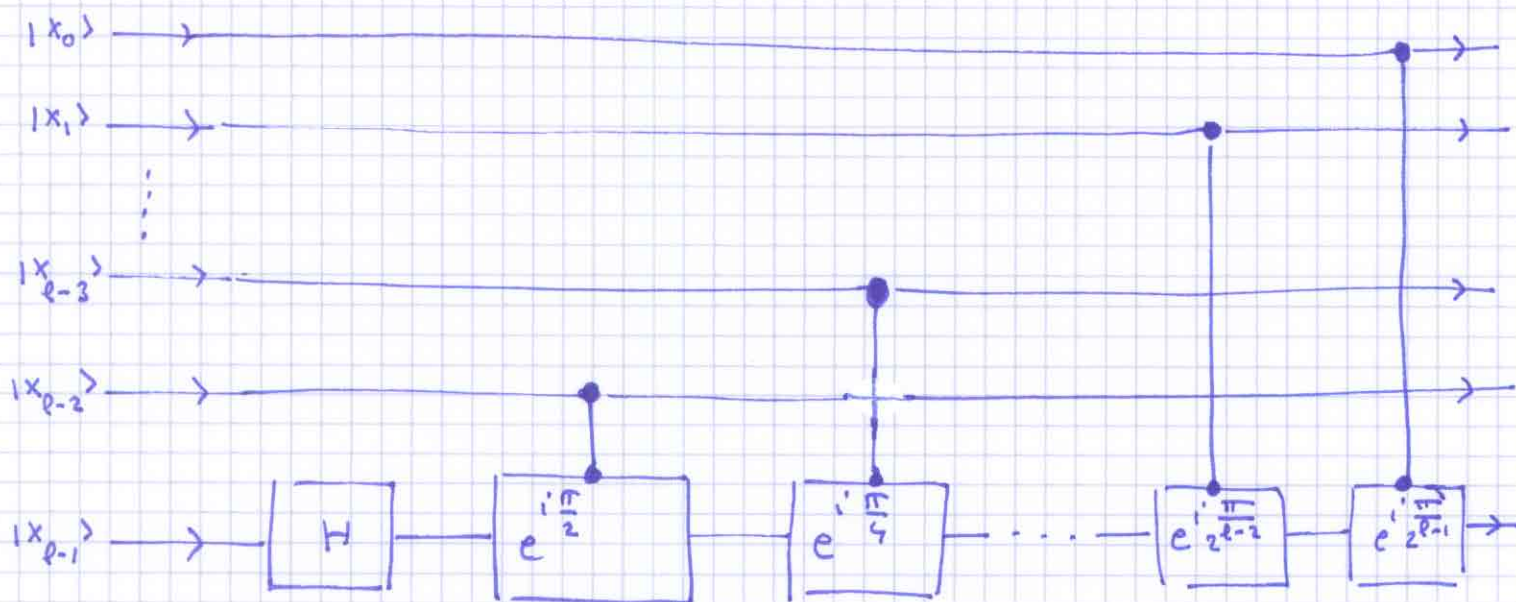
Let us look at the phase factor more closely. The binary expansion of x is :

$$x = 2^{m-1} \cdot x_{m-1} + \dots + 2^2 \cdot x_2 + 2^1 \cdot x_1 + 2^0 \cdot x_0.$$

and this implies that

$$e^{i \frac{\pi}{2^{l-1}} x} = (-1)^{x_{l-1}} \cdot e^{i \frac{\pi}{2} x_{l-2}} \cdot e^{i \frac{\pi}{4} x_{l-3}} \dots e^{i \frac{\pi}{2^{l-1}} x_0}.$$

So to obtain the l -th term in the product we may use the operations (Hadamard and double control phases).

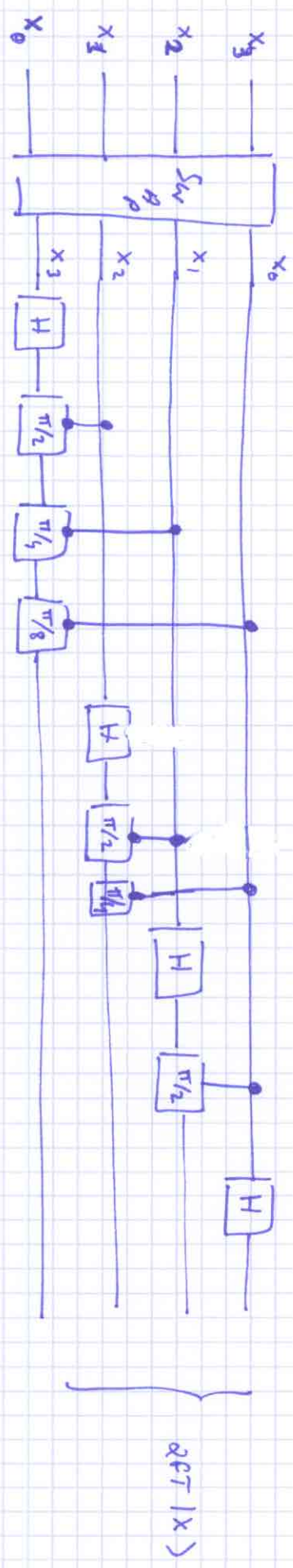


The output is $|x_0\rangle |x_1\rangle \dots |x_{l-2}\rangle \cdot \frac{1}{\sqrt{2}} \left(|0\rangle + e^{i \frac{\pi}{2^{l-1}} x} |1\rangle \right)$.

From these observations we deduce the full circuit for the QFT (Shor 1994)

$$\begin{cases} \text{input } |X\rangle = |x_{m-1} \dots x_0\rangle \\ \text{output } \text{QFT}|X\rangle = \frac{1}{2^{m/2}} \sum_{y=0}^{2^m-1} e^{2\pi i \frac{xy}{2^m}} |y\rangle. \end{cases}$$

For $m = 4$ qubits ($N = 2^4 = 16$) the circuit is (a 16×16 matrix):



To perform the SWAP here we can first swap x_0 & x_3 and then x_1 & x_2 . In any case the SWAP operation requires $\mathcal{O}(3M)$ CNOT gates and the rest of the QFT requires:

$$M + (M-1) + \dots + 1 = \frac{M(M-1)}{2} \text{ gates.}$$

H & phase gates H & phase gates H.

So the overall size of the QFT circuit is $\mathcal{O}(M^2) = \mathcal{O}(\ln N)^2$.

Remark: If one wants to work to finite accuracy

(which is the case in "practice") one can neglect the

phase gates $\frac{\pi}{2^k}$ for $\frac{1}{2^k} < \epsilon$ (ϵ a fixed accuracy).

Then the nb of gates becomes $O(n) = O(\ln N)$.

[The coefficient will be ϵ -dependent].